

# High- $J_c$ MgB<sub>2</sub> Josephson junctions with operating temperature up to 40 K

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Sandwich-type MgB<sub>2</sub>/MgO/MgB<sub>2</sub> Josephson junctions with Au or MgB<sub>2</sub> interconnection were fabricated using hybrid physical-chemical vapor deposited MgB<sub>2</sub> thin films and RF-magnetron-sputtered MgO barrier. The junctions show properties similar to those in high- $J_c$  Nb junctions with  $J_c$  up to 275 kA/cm<sup>2</sup> at 4 K, which remains nonzero up to 40 K. Critical current modulations by applied magnetic field and constant voltage steps under microwave radiation were observed. Combined with the larger energy gaps in MgB<sub>2</sub> than in Nb, the junctions presented here allow simple MgB<sub>2</sub> digital circuits to work over 20 K or with a clock speed above 1 THz. © 2010 American Institute of Physics. [doi:10.1063/1.3298366]

Superconducting integrated circuits using Josephson junctions have higher speed and lower power consumption than semiconductor circuits.<sup>1–3</sup> However, the only mature technology uses Nb (critical temperature  $T_c=9.25$  K), which requires a low operating temperature of about 4 K. Superconducting circuits using the 40 K superconductor MgB<sub>2</sub> can operate above 20 K (Ref. 4) cooled by more efficient, more compact, and less costly cryocoolers than those used at 4 K.<sup>5</sup> Such a breakthrough requires fabrication technologies for MgB<sub>2</sub> Josephson junctions with high critical current density  $J_c$  and multilayers consisting of superconducting, insulating and resistive films. There have been reports of Josephson junctions with one MgB<sub>2</sub> and one low- $T_c$  superconductor electrodes,<sup>6–8</sup> which cannot operate above the lower  $T_c$ . All-MgB<sub>2</sub> weak-link junctions<sup>9–13</sup> are not ideal for integrated circuits. All-MgB<sub>2</sub> sandwich-type Josephson junctions have been reported,<sup>14–16</sup> some with near-ideal Fraunhofer patterns, however, the  $J_c$  values are very low and drop quickly to zero at 20–25 K. In this letter, we report MgB<sub>2</sub>/MgO/MgB<sub>2</sub> sandwich-type Josephson junctions with high  $J_c$  up to 275 kA/cm<sup>2</sup> at 4 K, which remains nonzero up to 40 K.

The ultimate speed limit of Josephson junction-based rapid single flux quantum (RSFQ) circuits is proportional to the product  $I_c R_n$ , where  $I_c$  is the critical current and  $R_n$  the resistance of the junction in the normal state.<sup>1</sup> To eliminate hysteresis in  $I$ - $V$  characteristics, tunnel junctions in RSFQ circuits are usually critically damped by an external shunt resistor, and the maximum circuit speed is approximately proportional to  $J_c^{1/2}$ .<sup>17</sup> The highest operating speed for Nb circuit, 770 GHz, has been demonstrated by an RSFQ toggle flip-flop (TFF) frequency divider with  $J_c$  greater than 60 kA/cm<sup>2</sup>.<sup>18</sup> The energy gap in MgB<sub>2</sub> ( $\sim 2.2$  meV for the smaller of the two gaps<sup>19</sup>) is larger than that of Nb

( $\sim 1.5$  meV), implying a limiting clock speed for MgB<sub>2</sub> digital circuits of above 1 THz.<sup>2,18</sup>

The structure of the MgB<sub>2</sub>/MgO/MgB<sub>2</sub> sandwich-type Josephson junctions were illustrated in Fig. 1(a). A 100 nm-thick MgB<sub>2</sub> film was first grown on a SiC(0001) substrate by hybrid physical-chemical vapor deposition (HPCVD)<sup>20</sup> at 710 °C under conditions similar to those described in Ref. 8. The root-mean-square roughness of the MgB<sub>2</sub> film was 1.2–1.5 nm. The sample was then transferred in air to an rf magnetron sputtering system, where a 0.8–2 nm-thick MgO film was deposited at room temperature from a 2 inch MgO ceramic target at 25 W (130V) in 0.4 Pa pure Ar. The deposition rate was approximately 0.024 nm/min. The sample was

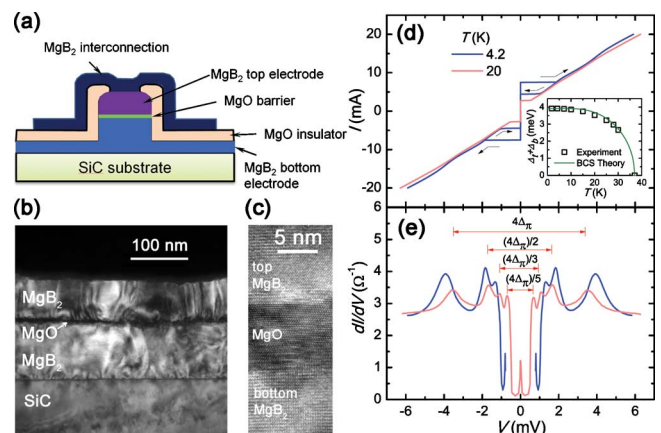


FIG. 1. (Color online) (a) Illustration of the junction structure. (b) Low magnification and (c) high resolution TEM cross-sectional view of the trilayer structure. The epitaxial relationship is MgB<sub>2</sub>[0001]||MgO[111]||MgB<sub>2</sub>[0001]||SiC[0001] and MgB<sub>2</sub>[1 $\bar{1}$ 00]|| $\times$ MgO[1 $\bar{1}$ 21]||MgB<sub>2</sub>[1 $\bar{1}$ 00]||SiC[1 $\bar{1}$ 00]. (d)  $I$ - $V$  and (e)  $dI/dV$ - $V$  characteristics of a junction at 4.2 K and 20 K. The arrows indicate the direction of current sweep at 4.2 K. (inset) The temperature dependence of the sum of energy gaps of both electrodes (symbols) compared to the Bardeen-Cooper-Schrieffer theory (line). Subgap peak positions for 20 K are indicated.

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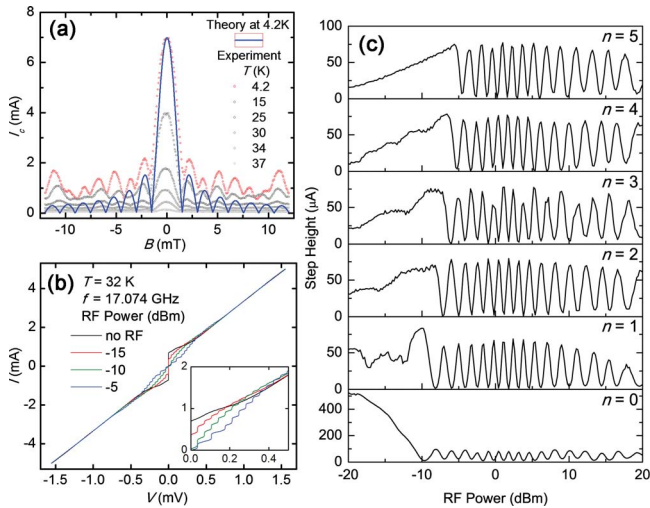


FIG. 2. (Color online) (a)  $I_c$  modulation by applied magnetic field at temperatures from 4.2 to 37 K (symbols) compared to the theoretical Fraunhofer pattern (line) and (b)  $I-V$  curves under 17.047 GHz microwave radiation for a  $20 \times 20 \mu\text{m}^2$  junction. (c) The microwave power dependence of the Shapiro steps for  $n$  from 0 to 5.

then transferred back to the HPCVD reactor for deposition of a 67 nm-thick top  $\text{MgB}_2$  electrode. Square-shaped junctions with sizes from  $4 \times 4$  to  $70 \times 70 \mu\text{m}^2$  were defined by photolithography and ion milling. After a subsequent sputter deposition of a 100 nm-thick MgO insulating layer and lift-off, an interconnection layer of gold or HPCVD  $\text{MgB}_2$  film was deposited. For genuine four-probe measurement, the contact to the top electrode was separated into two parts in the case of gold interconnection. Transmission electron microscope (TEM) cross-section images of a  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  trilayer are shown in Figs. 1(b) and 1(c). All layers are epitaxial with each other and with the SiC(0001) substrate. Both  $\text{MgB}_2/\text{MgO}$  interfaces are sharp without indication of interdiffusion.

The  $I-V$  characteristics of a  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  junction at 4.2 and 20 K are shown in Fig. 1(d). Hysteresis is pronounced at 4.2 K, becoming smaller as temperature increases, and eventually disappearing at around 20 K. The behavior can be accounted for by the resistively shunted junction model.<sup>1</sup> Figure 1(e) shows the  $dI/dV-V$  curves for the two temperatures. Two energy gaps exist in  $\text{MgB}_2$ , arising from the quasi-two-dimensional and  $ab$ -plane-confined  $\sigma$  bands [ $\Delta_\sigma(0)=7.4$  meV] and the three-dimensional  $\pi$  bands [ $\Delta_\pi(0)=2.2$  meV].<sup>21,22</sup> The two  $dI/dV$  peaks at around  $\pm 4$  mV correspond to the sum of the  $\pi$  gaps of the top and bottom  $\text{MgB}_2$  electrodes. Subgap peaks roughly at  $2\Delta_\pi/N$  ( $N=2, 3,$  and  $5$ ) are due to multiple Andreev reflection (MAR).<sup>23</sup> There is no observation of the  $\sigma$  gap because less than 1% of  $\sigma$ -band current is expected in  $c$ -axis tunneling.<sup>24</sup>

Figure 2(a) shows  $I_c$  of a  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  junction versus magnetic field  $B$  applied in the junction plane at temperatures from 4.2 to 37 K. A periodic modulation resembling the expected Fraunhofer pattern is observed. From the period of the oscillation the London penetration depth of the  $\text{MgB}_2$  film in  $c$ -axis is estimated to be 35 nm, consistent with theoretical prediction<sup>25</sup> and prior experimental measurements.<sup>26</sup> In contrast with the ideal Fraunhofer pattern, the minima of  $I_c(B)$  do not reach zero and the oscillation amplitude does not decay monotonically at high field, indicating a nonuniform current distribution over the junction area.

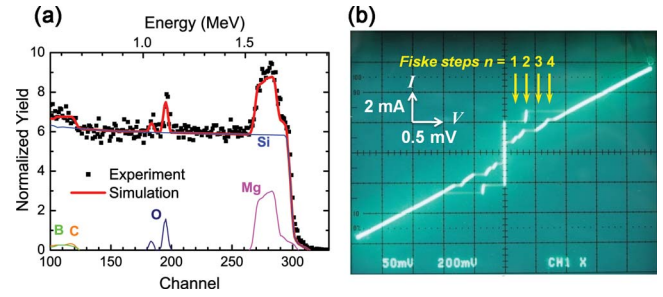


FIG. 3. (Color online) Measurement of the MgO barrier thickness by (a) RBS spectrum (squares, 3.05 MeV,  $\text{He}^{2+}$ ) and simulation (thick line) of an  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  trilayer, suggesting a well defined 1.6-nm-thick MgO layer sandwiched between the  $\text{MgB}_2$  electrodes, and by (b) the Fiske steps indicated by arrows on the  $I-V$  curve of a  $70 \times 70 \mu\text{m}^2$  junction under 0.1 mT field at 4.2 K with the MgO barrier thickness calculated to be about 1.2 nm.

This may partially be due to that the Josephson penetration depth of this junction ( $20 \mu\text{m}$ ) is comparable to the length of the junction ( $20 \mu\text{m}$ ) (Ref. 1) and possible conductive shorts within the barrier. Figure 2(b) shows the  $I-V$  curves of the junction under 17.1 GHz microwave radiations of various powers at 32 K. Shapiro steps were observed with step heights for different orders oscillating as the microwave power increases as shown in Fig. 2(c). The two experiments show unambiguously that the  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  device is a Josephson junction.

The TEM image in Fig. 1(c) suggests an MgO barrier thickness of about 5 nm for the trilayer, larger than 1.5 nm determined by the deposition time. The discrepancy may be due to the rippling of the MgO layer through the thickness of the TEM sample or the oxide layer formed during the sample transfer in air and plasma oxidation during MgO sputtering.<sup>27</sup> Figure 3(a) shows a Rutherford back scattering (RBS) spectrum and simulation for an  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  trilayer on SiC, the same as those used for the Josephson junctions. From the oxygen resonant scattering peak starting at channel number 200, the thickness of the MgO barrier is estimated to be 1.6 nm. Figure 3(b) is the  $I-V$  curve of a junction showing steps due to the Fiske resonance modes (marked by arrows),<sup>1</sup> which arise from the coupling between the Josephson current and the electromagnetic mode inside the dielectric barrier layer. Using  $\lambda_L=35$  nm and dielectric constant for MgO of 9.7,<sup>28</sup> the MgO layer thickness calculated from the steps is 1.2 nm, the same as determined from the deposition time. The two results are consistent with the thickness estimates from the deposition time, although errors can arise from oxygen vacancies in MgO for the RBS simulation and from the error in dielectric constant of MgO for the Fiske calculation.

Figure 4(a) shows an exponential dependence of  $J_c$  at 4.2 K on MgO barrier thickness, determined by the MgO deposition time. If we assume the simplest possible tunneling model for the Josephson current,  $J_c=J_0 \exp(-d/d_0)$ , where  $d_0=\hbar/(8mE_b)^{1/2}$  and  $E_b$  is the barrier height, we estimate  $E_b$  to be 0.8 eV, similar to that of MgO barriers in other tunnel junctions.<sup>29</sup> For an MgO layer thickness of 0.86 nm,  $J_c$  is  $275 \text{ kA}/\text{cm}^2$  at 4.2 K, comparable to the highest  $J_c$  achieved in Nb junctions.<sup>30</sup> The  $I_c(B)$  and  $I-V$  curve of this junction at 4.2 K are shown in Fig. 4(b). From the high  $J_c$  value, one infers an ultimate speed limit for a  $\text{MgB}_2$  RSFQ TFF frequency divider to be above 1 THz.<sup>2</sup>

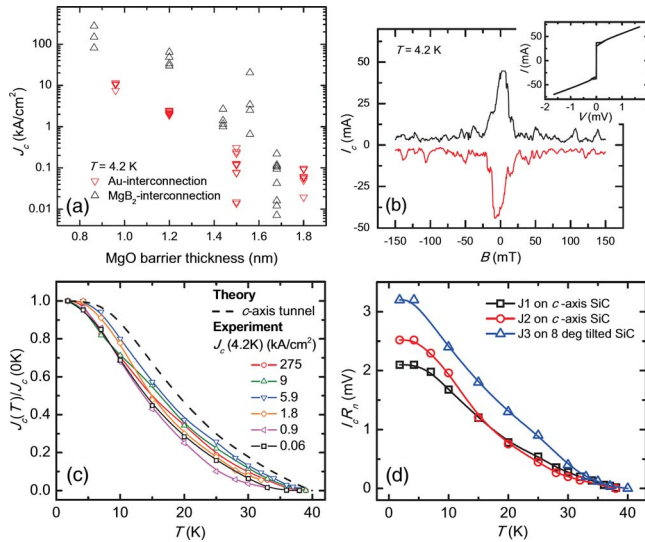


FIG. 4. (Color online) (a)  $J_c$  at 4.2 K for junctions of difference sizes on nine chips with Au or  $\text{MgB}_2$  interconnection and MgO barrier thickness varying from 0.8 to 1.8 nm. (b)  $I_c(B)$  and  $I$ - $V$  curve (inset) of a high  $J_c$  ( $275 \text{ kA/cm}^2$ ) junction ( $4 \times 4 \mu\text{m}^2$ ) at 4.2 K. (c) Normalized  $J_c(T)$  of various junctions (solid lines) compared to the theoretical calculation for  $\text{MgB}_2$   $c$ -axis Josephson tunnel junction (dashed line). (d)  $I_c R_n(T)$  of two  $c$ -axis (J1, J2) and one  $8^\circ$ -tilted (J3) junctions.

The behavior shown in Fig. 4(a) is different from that in Nb/Al- $\text{AlO}_x$ /Nb junctions, where the low- $J_c$  junctions are superconductor-insulator-superconductor tunnel junctions and the high- $J_c$  ( $>20 \text{ kA/cm}^2$ ) junctions are increasingly dominated by MAR. All the  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  junctions we have studied show properties similar to those in high- $J_c$  Nb/Al- $\text{AlO}_x$ /Nb junctions where MAR dominates. Whether some defect states or pinholes in the barriers are responsible for the MAR or not is a subject of ongoing investigation.

The temperature dependence of  $J_c$  for several  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  junctions is shown in Fig. 4(c) along with the result of a two-band tunneling theory.<sup>24</sup>  $J_c$  remains nonzero up to 40 K. The faster decrease in  $J_c$  with increasing temperature than the theory is possibly due to weak link contributions.

In Fig. 4(d),  $I_c R_n$  versus temperature is shown for three junctions. J3 was made on a SiC substrate with  $8^\circ$  off cut to expose  $ab$  planes of the  $\text{MgB}_2$  electrodes. The  $I_c R_n$  product at 4.2 K ranges from 2.1–3.1 mV, about half of the predicted value (4 and 6 mV for tunneling in  $c$  and  $ab$  directions, respectively).<sup>24</sup> J3 has the highest  $I_c R_n$ , possibly due to the  $\sigma$  band contribution.

For RSFQ circuits, superconducting interconnect is necessary.<sup>1</sup> Figure 4(a) shows that the  $J_c$  values for Au and  $\text{MgB}_2$  interconnections are comparable, indicating that the  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  junctions can survive the subsequent depositions of the insulator and  $\text{MgB}_2$  layers. The on-chip  $J_c$  spread of ten  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  junctions is  $\sim 10\%$ , which needs to be reduced to 1%–2% for RSFQ circuits.

In summary, we have fabricated  $\text{MgB}_2/\text{MgO}/\text{MgB}_2$  Josephson junctions with high  $J_c$  up to  $275 \text{ kA/cm}^2$  at 4 K, which remains nonzero up to 40 K. Their properties resemble high- $J_c$  Nb/Al- $\text{AlO}_x$ /Nb junctions. Improving the  $J_c$  uniformity, enhancing the  $I_c R_n$  product, and achieving nanometer-sized patterning for high- $J_c$  junctions ( $\sim 250 \text{ nm}$  for  $I_c$

$\sim 100 \mu\text{A}$ ) are important next steps. Nevertheless, the results presented here should enable simple  $\text{MgB}_2$  RSFQ circuits to demonstrate the ultrahigh speed of superconducting electronics at temperatures over 20 K.

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