High continuous wave power, 0.8 μ m-band, Al-free active-region diode lasers

J. K. Wade, ^{a)} L. J. Mawst, and D. Botez Reed Center for Photonics, University of Wisconsin-Madison 1415 Engineering Drive, Madison, Wisconsin 53706

M. Jansen, F. Fang, and R. F. Nabiev

Laser Group, Coherent, Inc. 2180 W. 190th Street, Torrance, California 90504

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Efficient, high-power, Al-free active-region diode lasers emitting at λ =0.83 μ m have been grown by low-pressure metalorganic chemical vapor deposition. Threshold-current densities as low as 220 A/cm², maximum continuous wave (cw) power of 4.6 W, and a maximum cw wallplug efficiency of 45% are achieved from 1 mm long, uncoated devices with In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P cladding layers. Further improvement is obtained by replacing the p-In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P cladding layer with thin (0.1 μ m) electron-blocking layers of Al_{0.85}Ga_{0.15}As and In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P, and a p-In_{0.5}(Ga_{0.9}Al_{0.1})_{0.5}P cladding layer. Such devices provide a record-high T_0 of 160 K and reach catastrophic optical mirror damage (COMD) at a record-high cw power of 4.7 W (both facets). The corresponding COMD power-density level (8.7 MW/cm²) is ~2 times the COMD power-density level for uncoated, 0.81- μ m-emitting AlGaAs-active devices. Therefore, 0.81- μ m-emitting, Al-free active-region devices are expected to operate reliably at roughly twice the power of AlGaAs-active region devices. © 1997 American Institute of Physics. [S0003-6951(97)04302-7]

High-power diode lasers emitting in the wavelength range $\lambda = 0.8 - 0.87 \mu m$ are of interest because of important applications such as pump sources for Nd:YAG lasers. Such lasers conventionally use AlGaAs in the active region for $\lambda \le 0.84 \mu m$, and thus suffer from short lifetimes and limited output powers by comparison to GaAs-active-region devices. The quaternary alloys InGa(As)P, lattice matched to GaAs, offer an attractive alternative to the AlGaAs/GaAs material system because: (1) the potential for better reliability, due to their inherent resistance to the formation of dark-line defects1 and lower surface recombination velocity of In-GaAsP compared to AlGaAs;² (2) a smaller increase in facet temperature with drive current, and (3) the potential for growing reliable diode lasers on Si substrates.³ However, the InGaAsP/GaAs material system has small conduction-band offsets, which cause diode lasers made of such material to suffer from massive carrier leakage. As a result, one obtains a relatively high threshold-current density, $J_{\rm th}$; 1,4,5 low internal efficiency, η_i ; 5,6 and low threshold-current characteristic temperature, T_0 . 5,7

A recent attempt to solve the problem of carrier leakage⁸ has been the use of high-band-gap material ($Al_{0.7}Ga_{0.3}As$) for the cladding layers which delivered high T_0 , promising reliability data, but also relatively high J_{th} . We have designed and fabricated high-power, efficient, Al-free active-region devices by employing two features: (high-band gap) cladding layers of $In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P$ and the "broad-waveguide" concept. The former significantly reduces carrier leakage, while the latter insures both a low internal loss coefficient, α_i (~ 2 cm⁻¹), since a large fraction of the optical energy is contained within the not intentionally doped (i.e., low loss) waveguide, as well as a large transverse spot size (0.5 μ m full-width at half-maximum). As a result, we report the lowest J_{th} (220 A/cm² for 1-mm-long devices) for 0.8 μ m band,

Al-free diode lasers as well as a record-high $T_0(160 \text{ K})$. Fur-

The structures are grown by low-pressure (50 mbar) metalorganic vapor phase epitaxy (LP-MOVPE) at a temperature of 700 °C (except for the p^+ -GaAs contact layer, which is grown at 650 °C to increase the Zn incorporation) on exactly oriented (100) n^+ -GaAs substrates. The source materials are trimethylgallium, trimethylindium, arsine, phosphine, and diethylzinc and silane for the p- and n-type dopants. The structures are evaluated by characterizing 100- μ m-wide stripe lasers formed by chemically etching through the p^+ -GaAs contact layer and defining metal-contact openings with a SiO2 mask. Devices of various lengths are tested under pulsed-current operation (2 kHz, 5 µs pulse width) to extract $J_{\rm th},~\alpha_i,~\eta_i,~{\rm and}~T_0.$ For cw measurements, devices are mounted junction-side down on diamond or Cu submounts and tested in a fixture employing thermoelectric cooling.

The first laser structure is shown in Fig. 1. The active region consists of a 150 Å $In_{0.04}Ga_{0.96}Al_{0.89}P_{0.11}$ quantum well surrounded by partially ordered 0.4 μ m $In_{0.5}Ga_{0.5}P$ optical confinement layers and 1.25 μ m $In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P$ cladding layers. A typical cw L-I curve is shown in Fig. 2 for a laser emitting at λ =0.83 μ m with uncoated facets, 100- μ m-wide stripe, and 1-mm-long cavity. The maximum output power is 4.6 W (both facets). The differential quantum efficiency, η_d , is 77% reflecting the low α_i (2 cm⁻¹) and high η_i (84%). Also shown in Fig. 2 is the wallplug efficiency, η_p , which reaches a maximum value of 45% at an output power of 1.23 W. This record-high η_p for Al-free 0.8- μ m-band diode lasers reflects low J_{th} , high η_d , small temperature sensitivity for J_{th} and η_d , and a relatively low

thermore, we have obtained from uncoated, 100- μ m-stripe devices record-high continuous wave (cw) output power (4.7 W) allowing the determination of the catastrophic-optical-mirror damage (COMD) power-density level for uncoated InGaAsP-active-layer devices in the 0.8 μ m band.

a)Electronic mail: wade@cae.wisc.edu

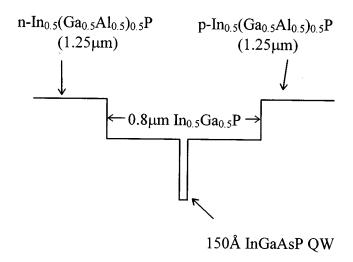


FIG. 1. Conduction-band edge of an Al-free active-region laser structure with 0.8- μ m-wide In_{0.5}Ga_{0.5}P waveguide layer and In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P cladding layers.

series resistance (0.2 Ω). Specifically, the characteristic temperature coefficients for $J_{\rm th}$ and η_d , T_0 and T_1 , 11 are 120 and 1220 K over the 20–60 °C temperature range. By comparison, the highest T_0 values reported for completely Alfree and 2-mm-long devices is 75 K for λ =0.875 μ m⁵ and 102 K for λ =0.805 μ m. Therefore, In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P cladding layers are highly effective at reducing carrier leakage with little penalty in electrical resistance.

For comparison, $J_{\rm th}$ vs inverse cavity length curves are shown for various Al-free 0.8- μ m-band devices in Fig. 3. As can be seen, the $J_{\rm th}$ values obtained for the structure in Fig. 1 are the lowest reported. Due to superior carrier confinement, $J_{\rm th}$ is only 365 A/cm 2 for 0.5-mm-long devices. The devices from Refs. 1 and 4 at λ =0.81 μ m, and Ref. 5 at 0.875 μ m are completely Al-free with correspondingly small conduction-band offsets (i.e., large carrier leakage) and subsequent high $J_{\rm th}$ values. As expected, the effect is particu-

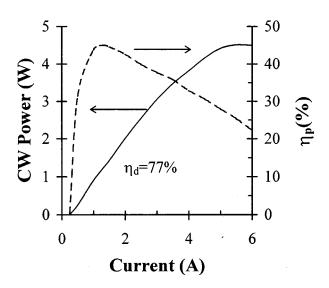


FIG. 2. cw L-I characteristics (both facets) and wallplug efficiency for a 1-mm-long laser with uncoated facets at 20 °C heatsink temperature using the structure shown in Fig. 1.

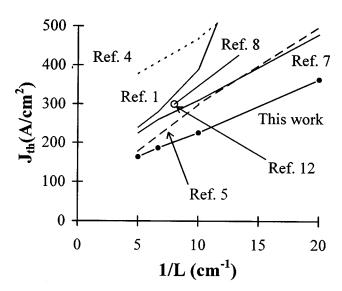


FIG. 3. Comparison between $J_{\rm th}$ vs inverse cavity length curves for the laser structure in Fig. 1 and other Al-free lasers in the 0.8 μ m-band (λ =0.81 μ m for Refs. 1, 4, 7, and 8; λ =0.83 μ m for Ref. 12; and λ =0.875 μ m for Ref. 5).

larly noticeable for devices operating at shorter wavelengths $(\lambda=0.81~\mu\text{m})^{1.4}$ with the notable exception of tensile-strained active-layer devices. At $\lambda=0.83~\mu\text{m}$ and for 1.25-mm-long devices, J_{th} for the structure in Fig. 1 is 65% of the lowest value reported for completely Al-free devices. With the addition of AlGaAs cladding layers [i.e., Ref. 8 ($\lambda=0.81~\mu\text{m}$)], devices show lower J_{th} and higher T_0 (~150 K for 1-mm-long devices) compared to Al-free 0.81 μ m devices. However, their J_{th} values are significantly higher than those for InGaAlP-clad devices.

To further lessen electron leakage, we studied the structure shown in Fig. 4, where the upper cladding layer of $p\text{-In}_{0.5}(\text{Ga}_{0.5}\text{Al}_{0.5})_{0.5}\text{P}$ has been replaced by thin (0.1 μ m) $p\text{-In}_{0.5}(\text{Ga}_{0.5}\text{Al}_{0.5})_{0.5}\text{P}$ and $p\text{-Al}_{0.85}\text{Ga}_{0.15}\text{As}$ layers, and a $p\text{-In}_{0.5}(\text{Ga}_{0.9}\text{Al}_{0.1})_{0.5}\text{P}$ layer for low thermal and electrical resistance. The results look promising as shown in Figs. 5(a) and 5(b) for 1.5-mm-long devices: T_0 has increased to 160 K

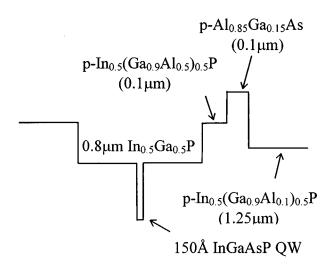


FIG. 4. Conduction-band edge of an Al-free active-layer laser structure with p-side Al_{0.85}Ga_{0.15}As/In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P carrier-blocking structure.

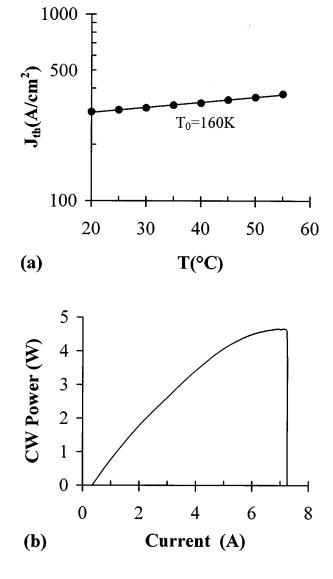


FIG. 5. (a) $J_{\rm th}$ vs temperature and (b) cw L-I characteristics (both facets) for a 1.5-mm-long laser with uncoated facets at 20 °C heatsink temperature made from the structure shown in Fig. 4.

and the maximum cw power achieved is 4.7 W (both facets). It would appear that $Al_{0.85}Ga_{0.15}As$ could not be an effective electron barrier because its band gap is less than that of $In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P$ (i.e., 2.11 eV vs 2.21 eV). However, the conduction-band offset ¹³ determines the electron-blocking capability. Recent studies ¹⁴ have shown that the conduction-band offset of $Al_{0.3}Ga_{0.7}As/In_{0.5}Ga_{0.5}P$ is the same as that for $In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P/In_{0.5}Ga_{0.5}P$ and increases with the Al content of $Al_xGa_{1-x}As$. That is, $Al_xGa_{1-x}As$ with x>0.3 will block electrons better than $In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P$ cladding-layer materials. This explains why the use of $Al_{0.7}Ga_{0.3}As$ as a cladding layer ⁸ and $Al_{0.85}Ga_{0.15}As$ as an electron blocking layer (this work) result in higher T_0 values than for $In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P$ -clad devices. The idea of using materials with advantageous conduction-band offsets for electron

blockage is not new: It has been proposed¹³ and implemented¹⁵ at long wavelengths (λ =1.55 μ m).

At 2.35 W per facet, the power is limited by COMD [Fig. 5(b)] at an internal power density of 8.7 MW/cm²; that is, 1.7 times that for GaAs-active-layer devices⁵ confirming that the surface recombination velocity of InGaAsP (λ =0.83 μ m) is lower than that of GaAs. AlGaAs used for 0.81 μ m emission has an even lower COMD than GaAs, ¹⁶ so that at λ =0.81 μ m InGaAsP-active layer devices are expected to have at least twice the COMD of AlGaAs devices and thus operate reliably at power levels twice as high.

In conclusion, 0.8- μ m-band Al-free active-region diode lasers with novel cladding layers have provided record-high cw output powers. Two laser structures were studied. The first used In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P cladding layers and provided 1-mm-long, uncoated devices with $J_{\rm th}=220{\rm A/cm^2}$, 4.6 W maximum cw output power and $\eta_{p,{\rm max}}=45\%$. The second used on the p-side 0.1- μ m-thick carrier-blocking layers of Al_{0.85}Ga_{0.15}As and In_{0.5}(Ga_{0.5}Al_{0.5})_{0.5}P followed by an In_{0.5}(Ga_{0.9}Al_{0.1})_{0.5}P cladding layer, and thus provided 4.7 W cw power and a COMD power density twice that of AlGaAsactive-layer devices. Therefore, these newly developed Alfree devices have the potential to be significantly more reliable than AlGaAs-based 0.81 μ m devices.

¹ D. Z. Garbuzov, N. Y. Antonishkis, A. D. Bondarev, A. B. Gulakov, S. Z. Zhigulin, N. I. Katsavets, A. V. Kochergin, and E. V. Rafailov, IEEE J. Quantum Electron. **QE-27**, 1531 (1991).

²E. Yablonovitch, R. Bhat, C. E. Zah, T. J. Gmitter, and M. A. Koza, Appl. Phys. Lett. **60**, 371 (1992).

³ T. Egawa, J. Dong, K. Matsumoto, T. Jimbo, and M. Umeno, IEEE Photonics Technol. Lett. 7, 1264 (1995).

⁴J. Diaz, I. Eliashevich, X. He, H. Yi, L. Wang, E. Kolev, D. Garbuzov, and M. Razeghi, Appl. Phys. Lett. 65, 1004 (1994).

⁵W. E. Plano, J. S. Major, Jr., and D. F. Welch, IEEE Photonics Technol. Lett. **6**, 465 (1994).

⁶H. J. Yi, J. Diaz, L. J. Wang, I. Eliashevich, S. Kim, R. Williams, M. Erdtmann, X. He, E. Kolev, and M. Razeghi, Appl. Phys. Lett. 66, 3251 (1995).

⁷ K. Uppal, A. Mathur, and P. D. Dapkus, IEEE Photonics Technol. Lett. 7, 1128 (1995).

⁸T. Fukunaga, M. Wada, H. Asano, and T. Hayakawa, Jpn. J. Appl. Phys. 34, L1175 (1995).

⁹I. B. Petrescu-Prahova, M. Buda, and T. G. van der Roer, IEICE Trans. Electron. **E77-C**, 1472 (1994).

¹⁰ D. Z. Garbuzov, J. H. Abeles, N. A. Morris, P. D. Gardner, A. R. Triano, M. G. Harvey, D. B. Gilbert, and J. C. Connolly, Proc. SPIE **2682**, 20 (1996).

¹¹L. J. Mawst, A. Bhattacharya, M. Nesnidal, J. Lopez, D. Botez, J. A. Morris, and P. Zory, Appl. Phys. Lett. 67, 2901 (1995).

¹² N. Y. Antonishkis, I. N. Arsent'ev, D. Z. Garbuzov, V. I. Kolyshkin, A. B. Komissarov, A. V. Kochergin, T. A. Nalet, and N. A. Strugov, Sov. Tech. Phys. Lett. **14**, 310 (1988).

¹³ R. F. Kazarinov and G. L. Belenky, IEEE J. Quantum Electron. QE-31, 423 (1995).

¹⁴I. Kim, Y. Cho, K. Kim, B. Choe, and H. Lim, Appl. Phys. Lett. **68**, 3488 (1996).

¹⁵ H. Murai, Y. Matsui, Y. Ogawa, and T. Kunii, Electron. Lett. 30, 2105 (1995).

¹⁶ K. Shigihara, Y. Nagai, S. Karakida, A. Takami, Y. Kokubo, H. Matsubara, S. Kakimoto, IEEE J. Quantum Electron. QE-27, 1537 (1991).